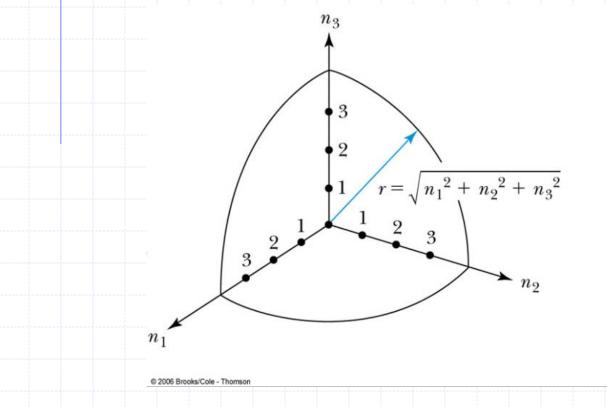
- As we mentioned previously, the Pauli exclusion principle plays an important role in the behavior of solids
 - The first quantum model we'll look at is called the free electron model or the Fermi gas model
 - It's simple but provides a reasonable description of the properties of the alkalai metals (Li, Na, ...) and the nobel metals (Cu, Ag, Au)

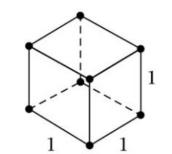
- Assume N electrons that are non-interacting (identical) fermions
- \blacktriangleright Assume they are confined to a box of V=L³
 - Hopefully you recall the energy levels of the 3D infinite potential well

$$E_{n_x n_y n_z} = \frac{\pi^2 \hbar^2}{2mL^2} \left(n_x^2 + n_y^2 + n_z^2 \right)$$

- Note as L grows large the discrete energy spectrum becomes continuous
- What is the ground state when N is large?
- Assume T=0 to start with

Consider a number space to help count the number of states





Cube of unit "volume"

Then we can write

$$E = \frac{\pi^2 \hbar^2}{2mL^2} \left(n_x^2 + n_y^2 + n_z^2 \right) \equiv E_1 r^2$$

The number of states up to radius r is just

$$N_r = \left(2\right) \left(\frac{1}{8}\right) \left(\frac{4}{3}\pi r^3\right)$$

We can rewrite this in terms of energy

$$N_r = \frac{1}{3}\pi \left(\frac{E}{E_1}\right)^{3/2}$$

 $N_r = \frac{1}{3}\pi \left(\frac{E_F}{E_1}\right)^{3/2}$

And at T = 0 the Fermi energy is the

highest occupied energy state

We can then solve for the Fermi energy

at

$$E_F = E_1 \left(\frac{3N}{\pi}\right)^{3/2} = \frac{\hbar^2 \pi^2}{2m} \left(\frac{3N}{\pi L^3}\right)^{2/3}$$

- Many metals have ~1 free electron/ion or n=N/L³ ~ 10²⁹ e/m³
- \succ Thus the Fermi energy E_F for most metals is ~ 10 eV
- > And the Fermi temperature $T_F = E_F/k_B$ is ~ 10⁵ K
 - Extremely high and much greater than T for room temperature ~ 300 K
 - But this is not the temperature of the electron gas, rather it is a measure of where the Fermi energy is

 \rightarrow N/L³

Selected Elements at $T = 300$ K						
Element	$\frac{N/V}{(imes 10^{28} { m m}^{-3})}$	Element	$\frac{N/V}{(\times 10^{28} \mathrm{m}^{-3})}$			
Cu	8.47	Mn (α)	16.5			
Ag	5.86	Zn	13.2			
Au	5.90	Cd	9.27			
Be	24.7	Hg (78 K)	8.65			
Mg	8.61	Al	18.1			
Ca	4.61	Ga	15.4			
Sr	3.55	In	11.5			
Ba	3.15	Sn	14.8			
Nb	5.56	Pb	13.2			
Fe	17.0					

 $\leftarrow \ge E_F \text{ and } T_F$

Table 9.4 Fermi Energies (T = 300 K), Fermi Temperatures, and Fermi Velocities for Selected Metals				
Element	$E_{\rm F}~({ m eV})$	$T_{\rm F}(\times10^4{\rm K})$	$u_{\rm F}$ (× 10 ⁶ m/s)	
Li	4.74	5.51	1.29	
Na	3.24	3.77	1.07	
K	2.12	2.46	0.86	
Rb	1.85	2.15	0.81	
Cs	1.59	1.84	0.75	
Cu	7.00	8.16	1.57	
Ag	5.49	6.38	1.39	
Au	5.53	6.42	1.40	
Be	14.3	16.6	2.25	
Mg	7.08	8.23	1.58	
Ca	4.69	5.44	1.28	
Sr	3.93	4.57	1.18	
Ba	3.64	4.23	1.13	
Nb	5.32	6.18	1.37	
Fe	11.1	13.0	1.98	
Mn	10.9	12.7	1.96	
Zn	9.47	11.0	1.83	
Cd	7.17	8.68	1.62	
Hg	7.13	8.29	1.58	
Al	11.7	13.6	2.03	
Ga	10.4	12.1	1.92	
In	8.63	10.0	1.74	
TI	8.15	9.46	1.69	
Sn	10.2	11.8	1.90	
РЬ	9.47	11.0	1.83	
Bi	9.90	11.5	1.87	
Sb	10.9	12.7	1.96	

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Comment

Once you know the Fermi energy you can also calculate the Fermi momentum and the Fermi velocity

$$p_F = \sqrt{2mE_F}$$

$$v_F = \frac{p_F}{m} = \sqrt{\frac{2E_F}{m}}$$

The Fermi velocity is the same order of magnitude as the orbital velocity of the outer electrons in an atom and ~10 times the mean thermal velocity of a non-degenerate electron gas

It's still non-relativistic however

The density of states g(E) is defined as

$$g(E) = \frac{dN}{dE} = \frac{\pi}{2} E_1^{-3/2} E^{1/2}$$

It can also be written in terms of E_F

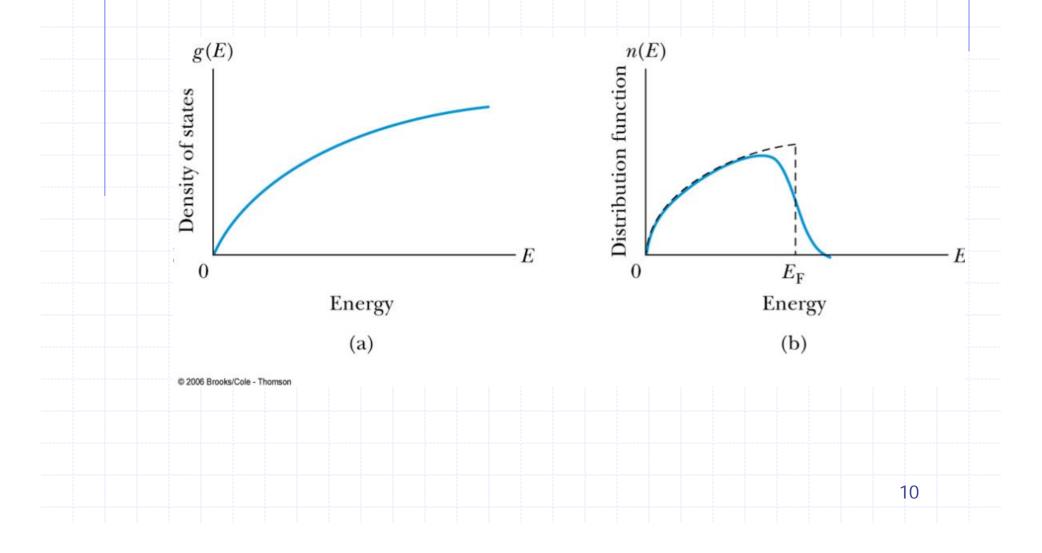
$$g(E) = \frac{3N}{2} E_F^{-3/2} E^{1/2}$$

Thus we can calculate n(E)=g(E)F_{FD}

Recall $F_{FD} = 1$ for $E < E_F$ and 0 for $E > E_F$

Then n(E) = g(E) for $E < E_F$ and 0 for $E > E_F$





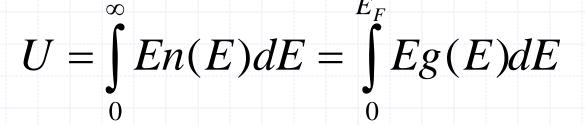
As we learned earlier, once you know n(E) you can calculate many thermodynamic quantities
 Sanity check

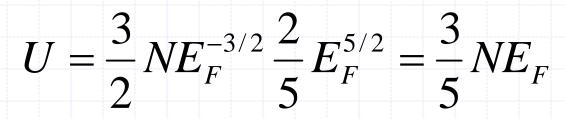
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$$N = \int_{0}^{\infty} n(E)dE = \int_{0}^{E_{F}} g(E)dE$$
$$N = \frac{3}{2}NE_{F}^{-3/2}\frac{2}{3}E_{F}^{3/2} = N$$







I used U here to indicate the total energy

Recall our discussion earlier in the semester on the molar heat capacity of solids

$$E = \frac{f}{2} N k_B T = \frac{f}{2} nRT$$
$$C_V = \frac{1}{n} \frac{dE}{dT} = \frac{f}{2} R$$

For solids $C_V = 3R \approx 25 \text{ J/molK}$

This is called the DuLong - Petit law

 But conduction electrons in a metal should contribute an additional (3/2)R
 The observed electronic contribution is only ~0.02R

We can calculate the electronic contribution in the free electron model

$$C_V = \frac{\partial U}{\partial T}$$

However up to this point we have been working at T=0 in order to calculate C_V so we have to calculate U(T) for T>0

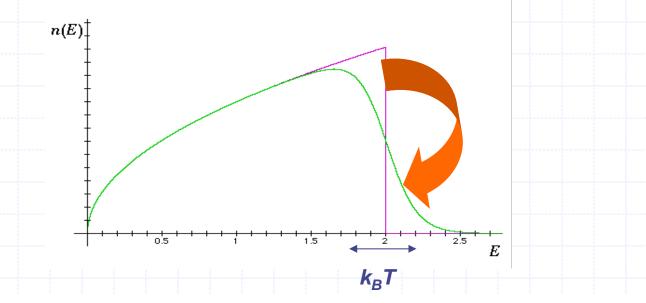
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$$U(T) = \int_{-\infty}^{\infty} En(E)dE = \int_{-\infty}^{\infty} Eg(E)F_{FD}dE$$

 $U(T) = \int_{0}^{\infty} Eg(E) \frac{1}{e^{\beta(E-E_F)} + 1}$

We could do the integral but it's more important to notice that in heating a Fermi gas, we populate some states above E_F and deplete some states below E_F

This modification is only important in a narrow energy range k_BT around E_F



The fraction of $k_{B}T/E_{F}$ electrons transferred to higher energy is $k_{R}T$ The energy increase of these electrons is So the increase in $N(k_{\rm B}T)^2/E_{\rm F}$ internal energy is > Making the molar heat $C_V = \frac{\partial U}{\partial T} \propto N \frac{k_B^2 T}{E_F} = R \frac{T}{T_F}$ capacity (for 1 mole, $N = N_{A_V}$)

The correct proportionality constant is found doing the integral

$$C_{V} = \frac{\pi^{2}}{2} N_{Av} k_{B} \frac{k_{B}T}{E_{F}} = \frac{\pi^{2}}{2} R \frac{T}{T_{F}}$$

- The point is, that at room temperature, T < < T_F and the electronic contribution to the heat capacity is small
- This is because only a small fraction of electrons around E_F can be thermally excited because of the Pauli exclusion principle
 - This was one of the big successes of the free electron model and not predicted classically

The free electron model for metals can also be used to calculate	
 Electrical conductivity Although to get agreement with experiment, 	
long collision lengths (mean free paths) were required	
 The source of how electrons can move ~ 1 centimeter (for a pure metal at low T) without scattering was unknown 	
Thermal conductivity	
 Although the classical calculation of thermal conductivity (fortuitously) gave good agreement with experiment also 	
18	3