We would like to move from the quantum theory of hydrogen to that for the rest of the periodic table One electron atom to multielectron atoms \succ This is complicated by the interaction of the electrons with each other and by the fact that the electrons are identical The Schrodinger equation for the two electron atom can only be solved by using approximation methods

- In classical mechanics, identical particles can be identified by their positions
 - In quantum mechanics, because of the uncertainty principle, identical particles are indistinguishable
 - This effect is connected with the Pauli exclusion principle and is of major importance in determining the properties of atoms, nuclei, and bulk matter

Consider two particles, then $\Psi(x,t) \rightarrow \Psi(x_1,x_2,t)$

The time dependent Schrodinger equation is

$$i\hbar \frac{\partial \Psi(x_1, x_2, t)}{\partial t} = \left[-\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x_1^2} - \frac{\hbar^2}{2m} \frac{\partial^2}{\partial x_2^2} + V(x_1, x_2) \right] \Psi(x_1, x_2, t)$$

The time independent Schrodinger equation is

$$-\frac{\hbar^2}{2m}\frac{\partial^2}{\partial x_1^2} - \frac{\hbar^2}{2m}\frac{\partial^2}{\partial x_2^2} + V(x_1, x_2)\bigg]\psi(x_1, x_2) = E\psi(x_1, x_2)$$

And note here we have used labels that identify particles 1 and 2

An important case is when the particles do not interact with each other

This is called the independent particle model

It's the starting point for the He atom e.g.

Then
$$V(x_1, x_2) = V_1(x_1) + V_2(x_2)$$

$$-\frac{\hbar^2}{2m}\frac{\partial^2}{\partial x_1^2} + V_1(x_1) - \frac{\hbar^2}{2m}\frac{\partial^2}{\partial x_2^2} + V_2(x_2)\psi(x_1, x_2) = E\psi(x_1, x_2)$$

This suggests we can write $\psi(x_1, x_2) = \psi(x_1)\psi(x_2)$ and $E = E_1 + E_2$ Which leads to

$$-\frac{\hbar^{2}}{2m}\frac{d^{2}}{dx_{1}^{2}} + V_{1}(x_{1})\bigg]\psi(x_{1}) = E_{1}\psi(x_{1})$$

$$-\frac{\hbar^{2}}{2m}\frac{d^{2}}{dx_{2}^{2}} + V_{2}(x_{2})\bigg]\psi(x_{2}) = E_{2}\psi(x_{2})$$
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Let's say particle 1 is in quantum state m and particle 2 is in quantum state n

$$\psi(x_1, x_2) = \psi_m(x_1)\psi_n(x_2)$$

If we interchange the two particles then

$$\psi(x_1, x_2) = \psi_m(x_2)\psi_n(x_1)$$

We get different wave functions (and probability densities) so the two particles are distinguishable

 Two particles are indistinguishable if we can exchange their labels without changing a measurable quantity

Thus neither of the solutions above is satisfactory

In fact there are two ways to construct indistinguishable wave functions

$$\psi_{s} = \frac{1}{\sqrt{2}} (\psi_{m}(x_{1})\psi_{n}(x_{2}) + \psi_{m}(x_{2})\psi_{n}(x_{1}))$$

$$\psi_{a} = \frac{1}{\sqrt{2}} (\psi_{m}(x_{1})\psi_{n}(x_{2}) - \psi_{m}(x_{2})\psi_{n}(x_{1}))$$

 Ψ_s is a symmetric wave function under particle interchange of 1↔2

- $\Psi_s \rightarrow \Psi_s$
- Ψ_a is a antisymmetric wave function under particle interchange of $1 \leftrightarrow 2$

• $\Psi_a \rightarrow -\Psi_a$

 The probability density remains unchanged under particle interchange of 1↔2

- All particles with integer spin are called bosons
 - Spin 0,1,2,...
 - Photon, pions, Z-boson, Higgs
 - All particles with half integer spin are called fermions
 - Spin 1/2, 3/2, ...
 - Electron, proton, neutron, quarks, ...
 - For the next few lectures we'll focus on the fermions (electrons)

- The wave function of a multi-particle system
 - of identical fermions is antisymmetric under interchange of any two fermions
 - of identical bosons is symmetric under interchange of any two bosons
 - The Pauli exclusion principle follows from these principles
 - No two identical electrons (fermions) can occupy the same quantum state

$$\psi_{a} = \frac{1}{\sqrt{2}} (\psi_{m}(x_{1})\psi_{n}(x_{2}) - \psi_{m}(x_{2})\psi_{n}(x_{1}))$$

for m = n (same quantum state)

$$\psi_{a} = \frac{1}{\sqrt{2}} (\psi_{m}(x_{1})\psi_{m}(x_{2}) - \psi_{m}(x_{2})\psi_{m}(x_{1})) = 0$$

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What about 3 particles?

> Symmetric

$$\psi_s = (\psi_m(x_1)\psi_n(x_2)\psi_o(x_3) + \psi_o(x_1)\psi_m(x_2)\psi_n(x_3) + \psi_n(x_1)\psi_o(x_2)\psi_m(x_3))$$

$$\psi_s = (\psi_m(x_1)\psi_n(x_2)\psi_o(x_3) + permutations)$$

Antisymmetric

• Use the Slater determinant $\psi_{a} = \frac{1}{\sqrt{3!}} \det \begin{vmatrix} \psi_{m}(1) & \psi_{m}(2) & \psi_{m}(3) \\ \psi_{n}(1) & \psi_{n}(2) & \psi_{n}(3) \\ \psi_{o}(1) & \psi_{o}(2) & \psi_{o}(3) \end{vmatrix}$ $\psi_{a} = \frac{1}{\sqrt{6}} (\psi_{m}(1)\psi_{n}(2)\psi_{o}(3) - \psi_{m}(1)\psi_{n}(3)\psi_{o}(2) + \psi_{m}(2)\psi_{n}(3)\psi_{o}(1) - \psi_{m}(2)\psi_{n}(1)\psi_{o}(3) + \psi_{m}(3)\psi_{n}(1)\psi_{o}(2) - \psi_{m}(3)\psi_{n}(2)\psi_{o}(1))$ 9

 \rightarrow Just a reminder, we are presently only working with the space wave functions We'll get to spin in a little bit \triangleright A consequence of identical particles is called exchange "forces" Symmetric space wave functions behave as if the particles attract one another Antisymmetric wave functions behave as if the particles repel one another



Let's return to the infinite well problem only now with two particles

Case 1 distinguishable particles

$$\psi(x_1, x_2) = \psi_m(x_1)\psi_n(x_2)$$

Case 2 identical particles (symmetric)

$$\psi(x_1, x_2) = \frac{1}{\sqrt{2}} (\psi_m(x_1) \psi_n(x_2) + \psi_m(x_2) \psi_n(x_1))$$

Case 3 identical particles (antisymmetric)

$$\psi(x_1, x_2) = \frac{1}{\sqrt{2}} (\psi_m(x_1)\psi_n(x_2) - \psi_m(x_2)\psi_n(x_1))$$

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What is the probability that both particles will be on the left side of the well?

$$P = \int_{0}^{L/2} \int_{0}^{L/2} |\psi(x_1, x_2)| dx_1 dx_2$$

Case 1

$$P = \int_{0}^{L/2} |\psi_m(x_1)|^2 dx_1 \int_{0}^{L/2} |\psi_n(x_2)|^2 dx_2$$

$$P = \frac{2}{L} \frac{2}{L} \int_{0}^{L/2} \frac{\sin^2 m \pi x}{L} dx \int_{0}^{L/2} \frac{\sin^2 n \pi x}{L} dx$$
$$P = \frac{4}{L^2} \frac{L}{4} \frac{L}{4} = \frac{1}{4}$$

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are more likely on opposite sides (repel) 14

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 - of identical fermions is antisymmetric under interchange of any two fermions
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 - The Pauli exclusion principle follows from these principles
 - No two identical electrons (fermions) can occupy the same quantum state

$$\psi_{a} = \frac{1}{\sqrt{2}} (\psi_{m}(x_{1})\psi_{n}(x_{2}) - \psi_{m}(x_{2})\psi_{n}(x_{1}))$$

for m = n (same quantum state)

$$\psi_{a} = \frac{1}{\sqrt{2}} (\psi_{m}(x_{1})\psi_{m}(x_{2}) - \psi_{m}(x_{2})\psi_{m}(x_{1})) = 0$$

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- The Pauli exclusion principle is the basis of the periodic table
 - This is why all the electron's don't simply fall into the ground state – because they are fermions
 - Recall there were 4 quantum numbers that specified the complete hydrogen wave function: n, l, m_l, and m_s
 - No two electrons in the same atom can have these same quantum numbers





Ionization Potential



Hydrogen (H) 1S¹ ≻Helium (He) 1s² – closed shell and chemically inert ≻Lithium (Li) Is²2s – valence +1, low I, partially screened, chemically very active ➢ Beryllium (Be) Is²2s² – Closed subshell but the 2s electrons can extend far from the nucleus



Fluorine (F)
 1s²2s²2p⁵ – Very chemically active because it can accept an electron to become a closed shell

≻Neon (Ne)

1s²2s²2p⁶ – Like He

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Gr	roups:	1	2											13	14	15	16	17	18	
		¹ н																	2 He	
		24																	142	}
		3	4	1										5	6	7	8	9	10	
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-	$4d^{10}5s^25p^6$	Cs	Ba	La	Hf	Ta	w	Re	Os	Ir	Pt	Au	Hg	TI	Pb	Bi	Po	At	Rn	
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- Periodic table is arranged into groups and periods
- ➢ Groups
 - Have similar shell structure hence have similar chemical and physical properties
 - Examples are alkalis, alkali earths, halogens, inert gases
- Periods
 - Correspond to filling d and f subshells
 - Examples are transition metals (3d,4d,5d), lanthanide (4f), and actinide (5f) series
 - Because there are many unpaired electrons, spin is important for these elements and there are large magnetic effects