

Finite Square Well

\rightarrow In regions I and III

$$\frac{-\frac{\hbar^2}{2m}\frac{d^2\psi_I}{dx^2} + V\psi_I = E\psi_I}{\frac{d^2\psi_I}{dx^2} - k_I^2\psi_I = 0 \text{ where } k_I = k_{III} = \frac{\sqrt{2m(V_0 - E)}}{\hbar}$$

➢ In region II

$$-\frac{\hbar^2}{2m}\frac{d^2\psi_{II}}{dx^2} = E\psi_{II}$$

$$\frac{d^2\psi_{II}}{dx^2} + k_{II}^2\psi_{II} = 0 \text{ where } k_{II} = k = \frac{\sqrt{2mE}}{\hbar}$$

Finite Square Well

Solutions in region I and III

$$Ae^{k_Ix} + Be^{-k_Ix} \rightarrow Ae^{k_Ix}$$

$$Ee^{k_{III}x} + Fe^{-k_{III}x} \rightarrow Fe^{-k_{III}x}$$

Solutions in region II

$$Ce^{ikx} + De^{-ikx}$$
 or

$C\sin kx + D\cos kx$



Finite Square Well

- Unfortunately we will not go much further
 - As with the infinite well, application of the boundary conditions leads to energy quantization
 - Although there are 4 equations for 4 unknowns the energy levels must be found numerically or graphically
 - As with the infinite well, the n'th eigenfunction will have n-1 nodes



- The (simple) harmonic oscillator is one of the most important physical systems in physics
 - Any physical system in the neighborhood of a stable equilibrium position can be approximated by a harmonic oscillator (in the limit of small oscillations)
 - Vibrations of atoms in a molecule, oscillations of atoms in a crystal, ...

Review from classical physics Mass on a spring $F_x = -kx$ (Hooke's law) $V = -W = -\int F dx = \frac{1}{2}kx^2$ $m\frac{d^2x}{dt^2} = -kx$ has solution $x = A\cos(\omega t - \phi)$ where $\omega = \sqrt{\frac{k}{m}}$ $E = T + V = \frac{p^2}{2m} + \frac{1}{2}kx^2$ $=\frac{m^2A^2\omega^2}{2m}\sin^2(\omega t-\phi)+\frac{1}{2}kA^2\cos^2(\omega t-\phi)$ $=\frac{1}{2}m\omega^2 A^2$ 8

Consider the neighborhood of the minimum of an arbitrary potential

Expand V(x) in a Taylor's series about the minimum at $x = x_0$

$$V(x) = a + b(x - x_0) + c(x - x_0)^2 + \dots$$

where

$$a = V(x_0) \equiv 0$$
$$b = \frac{dV}{dx}|_{x=x_0} = 0$$
$$c = \frac{1}{2} \left(\frac{d^2 V}{dx^2}\right)|_{x=x_0}$$

()

then

$$V(x) = \frac{1}{2}c(x - x_0)^2$$



\rightarrow We want to solve

$$-\frac{\hbar^2}{2m}\frac{d^2\psi}{dx^2} + \frac{1}{2}kx^2\psi = E\psi$$

Unfortunately this is not trivial. And an "easier" method (using ladder operators) is beyond the scope of the class.

We will obtain the asymptotic solution however

We start with the time independent Schrodinger equation

$$-\frac{\hbar^2}{2m}\frac{d^2\psi}{dx^2} + \frac{1}{2}kx^2\psi = E\psi$$

let
$$\omega = \sqrt{\frac{k}{m}}, \varepsilon = \frac{2E}{\hbar\omega}, y = \sqrt{\frac{m\omega}{\hbar}}x$$

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then
$$\frac{d^2\psi}{dy^2} + (\varepsilon - y^2)\psi = 0$$



Since the asymptotic solution and its derivative must vanish at infinity (why?) it must be that

$$\frac{d\psi_{\infty}}{dy} = \pm y\psi_{\infty}$$

C=0

The solution at infinity is thus

$$\psi_{\infty}(y) = e^{-\frac{y^2}{2}}$$

And that was easy compared to the full solution



The normalized wave functions = ψ(x) are given in Thornton and Rex (p223) and are not so illuminating

You should have be able to sketch the first few wave functions based on your experience with the infinite and finite wells and your knowledge of ψ in general



















In solving the harmonic oscillator problem for all y, the requirement that the wave function vanish at infinity leads to quantized energy

The energy levels of the harmonic oscillator are

$$E = \left(n + \frac{1}{2}\right)\hbar\omega \text{ with } n = 0, 1, 2, \dots$$

• The zero point energy is $E_0 = h\omega/2$

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The energy levels are equally spaced

This is similar to the Planck relation for radiation

 Recall we said that standing EM waves in the blackbody cavity arose from harmonic oscillators in the walls

← **Problem**

Use the Heisenberg uncertainty principle to estimate the ground state energy of the harmonic oscillator

← > Problem

Let ψ(x,0)=A{ψ₀(x) + ψ₁(x)} for the harmonic oscillator

- Normalize ψ
- Find |ψ(x,t)|²

← ▶ Problem

- Consider a potential that has
 - V(x)= ∞ for x < 0
 - $V(x) = 1/2kx^2$ for x > 0
- What are the possible energy levels?

 Einstein used the quantum solution of the harmonic oscillator to solve the temperature dependence of the specific heats of solids
 Classically

$$E = \frac{f}{2} Nk_B T = \frac{f}{2} nRT$$
$$C_V = \frac{1}{n} \frac{dE}{dT} = \frac{f}{2} R$$

For solids $C_V = 3R \approx 25 \text{ J/molK}$

This is called the DuLong - Petit law



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➢ If atoms are modeled by quantum harmonic oscillators, explain why specific heats drop as T→0