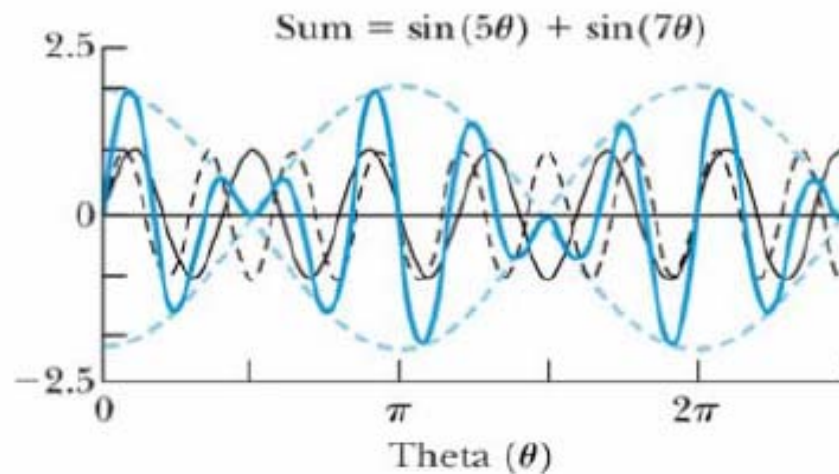


# Waves

- The first term is the wave and the second term is the envelope

$$\Psi(x, t) = 2A \cos(k_{av}x - \omega_{av}t) \cos\left(\frac{\Delta k}{2}x - \frac{\Delta\omega}{2}t\right)$$



# Waves

- What is the size in space and time of the wave packet?
- For the moment, let's define the size of the localized wave to be  $\Delta x = x_2 - x_1$  where points 1 and 2 are points where the envelope is zero

$$\frac{\Delta k}{2} x_2 - \frac{\Delta k}{2} x_1 = \pi$$

$$\Delta k \Delta x = 2\pi$$

and similarly

$$\frac{\Delta t}{2} \omega_2 - \frac{\Delta t}{2} \omega_1 = \pi$$

$$\Delta \omega \Delta t = 2\pi$$

# Waves

- We learned a lot from just summing two waves, but to localize the particle and remove the auxiliary waves we must sum/integrate over many waves
- This can be written as a Fourier integral

$$\Psi(x, t) = \int_{-\infty}^{\infty} dk A(k) e^{i(kx - \omega t)}$$

- The above form is equivalent to the one used in your book

$$\Psi(x, t) = \int_{-\infty}^{\infty} dk A(k) \cos(kx - \omega t)$$

# Waves

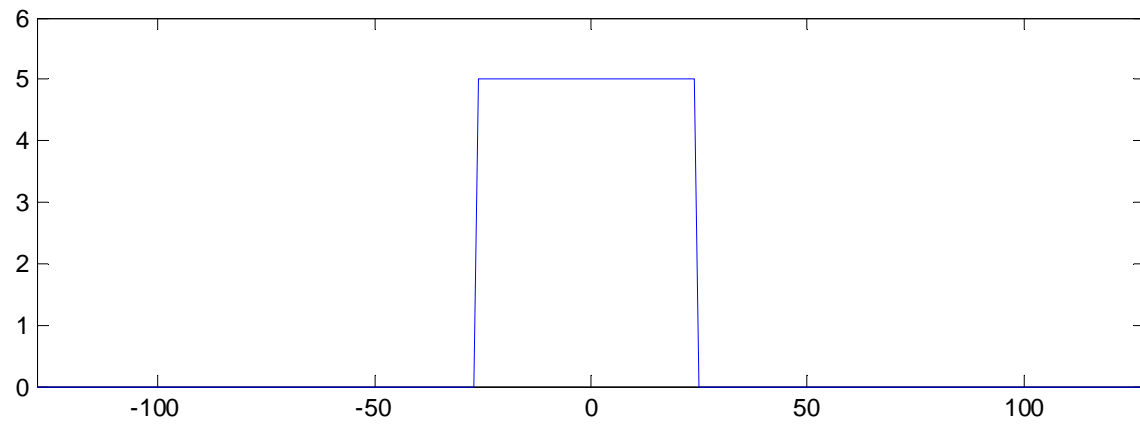
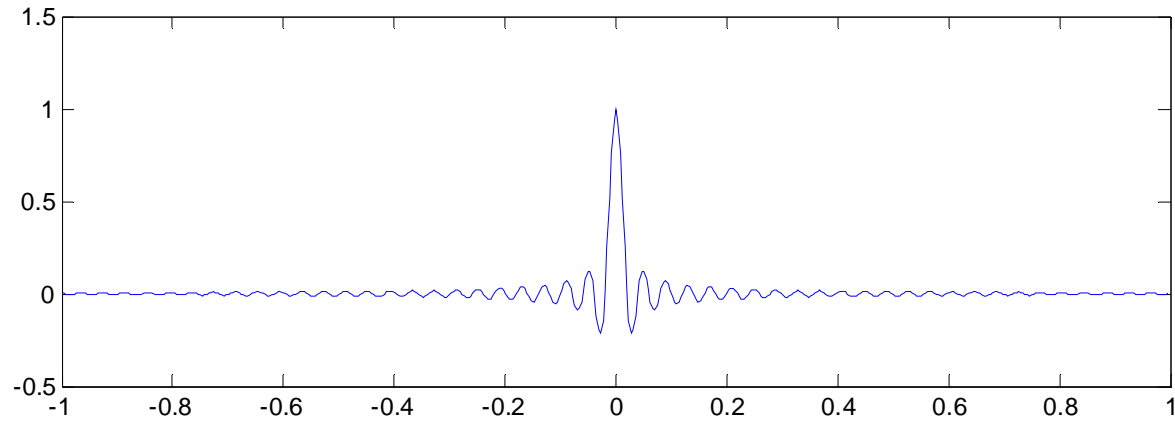
## ➤ Problem 5.31

$$\begin{aligned} \text{let } \tilde{A}(k) &= A_0 \text{ for } k = k_0 \mp \Delta k \\ &= 0 \text{ otherwise} \end{aligned}$$

Sketch the envelope term, the oscillating term,  
and  $|\Psi(x,0)|^2$

Defining  $\Delta x$  to be the FWHM, find  $\Delta x \Delta k$

# Waves



# Waves

➤ These  $\Delta x \Delta k$  relations mean

- If you want to localize the wave to a small position  $\Delta x$  you need a large range of wave numbers  $\Delta k$
- If you want to localize the wave to a small time domain  $\Delta t$  you need a large range of frequencies (bandwidth)  $\Delta \omega$

# Waves

➤ Consider the wave packet at  $t=0$

$$\Psi(x,0) = \int_{-\infty}^{\infty} dk A(k) e^{ikx}$$

then  $A(k)$  is the Fourier transform of  $\Psi(x,0)$

$$A(k) = \int dx \Psi(x,0) e^{-ikx}$$

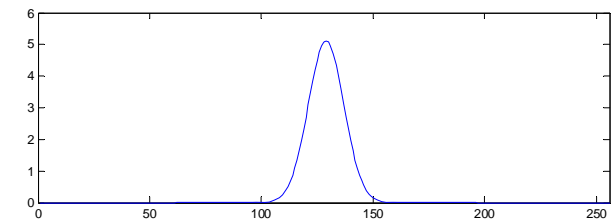
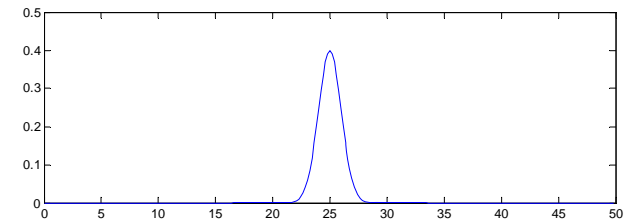
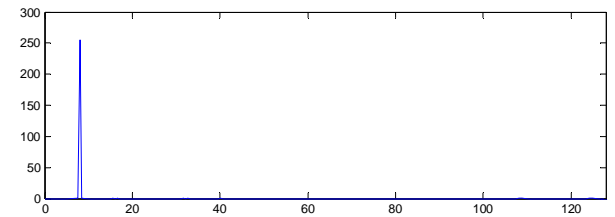
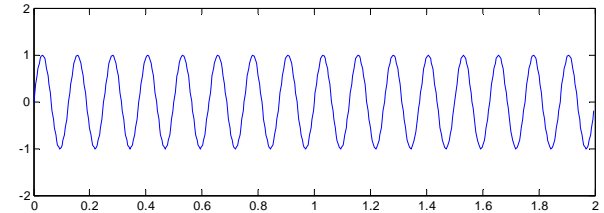
➤ The Fourier transform just transforms one function into another

- Position space is transformed into wave number space
- Time space is transformed into frequency space

# Waves

## ➤ Some simple examples

- Space (cosine)
- Wave number (delta function)
- Space (gaussian)
- Wave number (gaussian)





# Waves

➤ A particularly useful wave packet is the gaussian wave packet

- Both the wave packet and the Fourier transform are gaussian functions
- It's useful because the math is easier

$$\Psi(x,0) = \int_{-\infty}^{\infty} dk A(k) e^{ikx}$$

$$\text{with } A(k) = e^{-\alpha(k-k_0)^2/2}$$

- $A(k)$  is centered at  $k_0$  and falls away rapidly from the center

# Waves

- Doing the integral (by first changing variables to  $q=k-k_0$ ) gives

$$\Psi(x,0) = \int_{-\infty}^{\infty} dk A(k) e^{ikx}$$

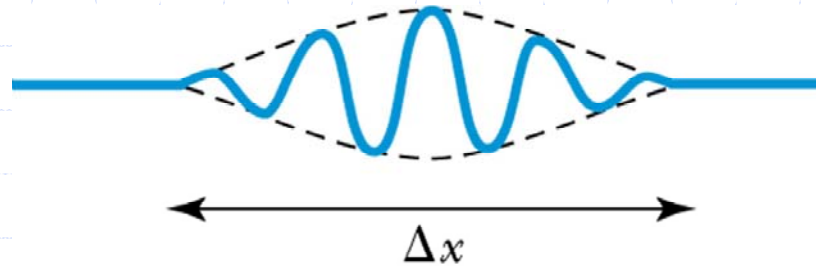
with  $A(k) = e^{-\alpha(k-k_0)^2/2}$  gives

$$\Psi(x,0) = \sqrt{\frac{2\pi}{\alpha}} e^{ik_0x} e^{-\frac{x^2}{2\alpha}}$$

- The factors represent an oscillating wave with wave number  $k_0$  and a modulating envelope
  - The wave function  $\Psi(x,0)$  and its Fourier transform  $A(k)$  are both gaussians

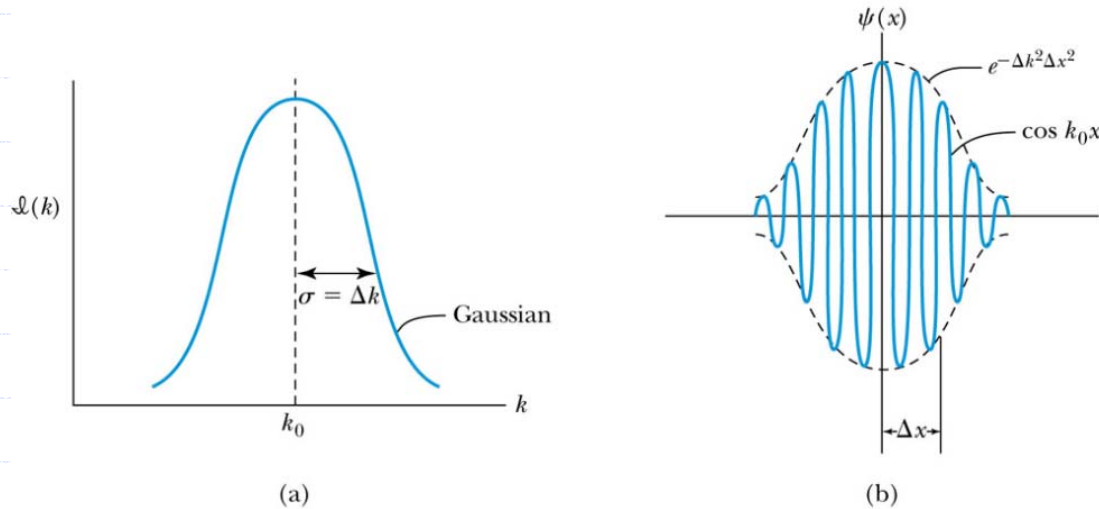
# Waves

➤ An idealized wave packet



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➤ A gaussian wave packet



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# Waves

➤ We can define the width  $\Delta x$  of the gaussian to be the width between which the gaussian function is reduced to  $1/\sqrt{e}$

$$e^{-\frac{1}{2}} = \frac{1}{\sqrt{e}}$$

$$-\frac{1}{2} = -\frac{x^2}{2\alpha} \text{ or } \Delta x = 2\sqrt{\alpha}$$

and doing the same for the k - space gaussian

$$-\frac{1}{2} = -\alpha \frac{(\Delta k)^2}{2} \text{ or } \Delta k = \frac{2}{\sqrt{\alpha}}$$

$$\text{then } \Delta x \Delta k = 4$$

# Waves

- Slightly different definitions of the width would give slightly different results
  - Still, the more localized the wave packet, the larger the spread of associated wave numbers
- A general result of Fourier integrals is

$$\Delta x \Delta k > \frac{1}{2} \text{ and since } p = \hbar k$$

$$\Delta x \Delta p > \frac{\hbar}{2}$$

- This is the Heisenberg uncertainty principle
  - This is a QM result arising from a property of Fourier transforms and de Broglie waves

# Waves

## ➤ Time evolution of the wave packet

The wave function at  $t = 0$  is

$$\Psi(x,0) = \int_{-\infty}^{\infty} dk A(k) e^{ikx}$$

The wave function at later times is

$$\Psi(x,t) = \int_{-\infty}^{\infty} dk A(k) e^{i(kx - \omega t)}$$

One can show (and we won't)

$$\Delta x(t) = \sqrt{(\Delta x(0))^2 + \left( \frac{\hbar t}{2m\Delta x(0)} \right)^2}$$

This means the wave is dispersing with time

# Waves

➤ Consider an electron confined to 0.1nm with  $k_0=0$  (at rest). When will the wave packet expand to  $\sqrt{2}$  times its initial size

$$\Delta x^2(t) = \left( \frac{\hbar t}{2m\Delta x(0)} \right)^2$$

$$t = \frac{(\Delta x)^2 2m}{\hbar}$$

$$t = \frac{(0.1nm)^2 (2)(0.511 \times 10^6 eV)}{(197.33 eVnm)(3 \times 10^8 m/s) \left( 10^9 \frac{nm}{m} \right)}$$

$$t = 1.7 \times 10^{-16} s$$