

# Waves

➤ Electrons (and other particles) exhibit both particle and wave properties

- Particles are described by classical or relativistic mechanics
- Waves will be described by wave mechanics = quantum mechanics

# Waves

➤ Something disturbing we'll have to address

deBroglie relations  $p = \frac{h}{\lambda}$  and  $E = hf$

wave velocity  $v = \lambda f$

$$\text{so } v = \frac{h}{p} \frac{E}{h} = \frac{E}{p} = \frac{mV^2 / 2}{mV} = \frac{V}{2}$$

➤ The wave and particle seem to be moving at different velocities

# Waves

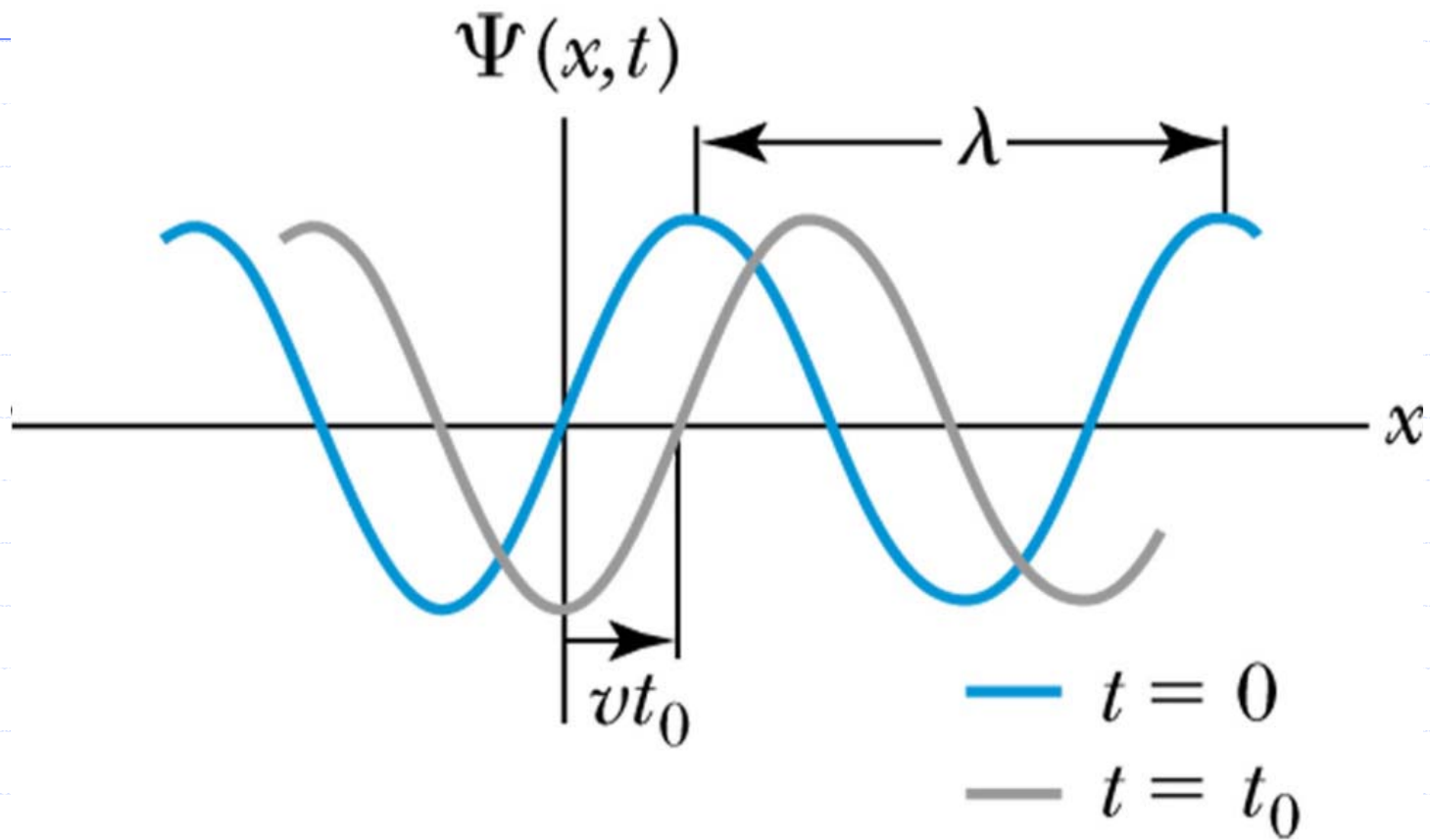
- We'll use  $\Psi(x,t)$  to represent the wave function of a particle (electron)
- In lecture 1, we wrote down the **classical** wave equation that describes a propagating disturbance

$$\frac{d^2\Psi}{dx^2} = \frac{1}{v^2} \frac{d^2\Psi}{dt^2}$$

- A solution to the classical wave equation is

$$\Psi(x,t) = A \sin\left[\frac{2\pi}{\lambda}(x - vt)\right]$$

# Waves



# Waves

$$\frac{d\Psi}{dx} = A \frac{2\pi}{\lambda} \cos\left[\frac{2\pi}{\lambda}(x - vt)\right]$$

$$\frac{d^2\Psi}{dx^2} = -A \left(\frac{2\pi}{\lambda}\right)^2 \sin\left[\frac{2\pi}{\lambda}(x - vt)\right]$$

$$\frac{d\Psi}{dt} = -vA \frac{2\pi}{\lambda} \cos\left[\frac{2\pi}{\lambda}(x - vt)\right]$$

$$\frac{d^2\Psi}{dt^2} = -Av^2 \left(\frac{2\pi}{\lambda}\right)^2 \sin\left[\frac{2\pi}{\lambda}(x - vt)\right]$$

$$\text{so } \frac{d^2\Psi}{dx^2} = \frac{1}{v^2} \frac{d^2\Psi}{dt^2} \text{ is satisfied}$$

# Waves

➤ We normally write

$$\Psi(x, t) = A \sin \left[ \frac{2\pi}{\lambda} (x - vt) \right]$$

as

$$\Psi(x, t) = A \sin(kx - \omega t)$$

where

$$k = \frac{2\pi}{\lambda} \text{ and } \omega = 2\pi f = \frac{2\pi}{T}$$

# Waves

➤ More generally we write

$$A \sin(kx - \omega t + \phi)$$

➤ Thus for  $\phi = \pi/2$  we have

$$A \cos(kx - \omega t)$$

➤ The phase velocity is defined as the velocity of a point on wave with a given phase (e.g the velocity of the peak of the wave)

$$v_{phase} = \lambda f = \frac{\lambda}{T} = \frac{\omega}{k}$$

# Waves

➤ The general harmonic wave can also be written as

$$A \sin(kx - \omega t + \phi) = B \sin(kx - \omega t) + C \cos(kx - \omega t)$$

where  $B = \cos \phi$  and  $C = \sin \phi$



# Waves

Using Euler's formula

$$e^{i\theta} = \cos \theta + i \sin \theta$$

we can write

$$\cos \theta = \frac{1}{2} (e^{i\theta} + e^{-i\theta})$$

$$\sin \theta = \frac{1}{2i} (e^{i\theta} - e^{-i\theta})$$

Then other solutions to the wave equation are

$$\text{Im}\{A \exp[i(kx - \omega t + \phi)]\} \text{ or } \text{Re}\{\dots\}$$

$$\frac{1}{2i} A \exp[i(kx - \omega t + \phi)] + \text{complex conjugate (c.c.)}$$

# Waves

## ➤ Notation

We often write these solutions without the  $\text{Re}$ ,  $\text{Im}$ ,  $\frac{1}{2}$ , or  $\frac{1}{2i}$

$\text{Im}\{A \exp[i(kx - \omega t + \phi)]\}$  becomes just

$$A \exp[i(kx - \omega t + \phi)]$$

and furthermore

$$A \exp[i(kx - \omega t + \phi)] = A \exp[i\phi] \exp[i(kx - \omega t)]$$

can be written

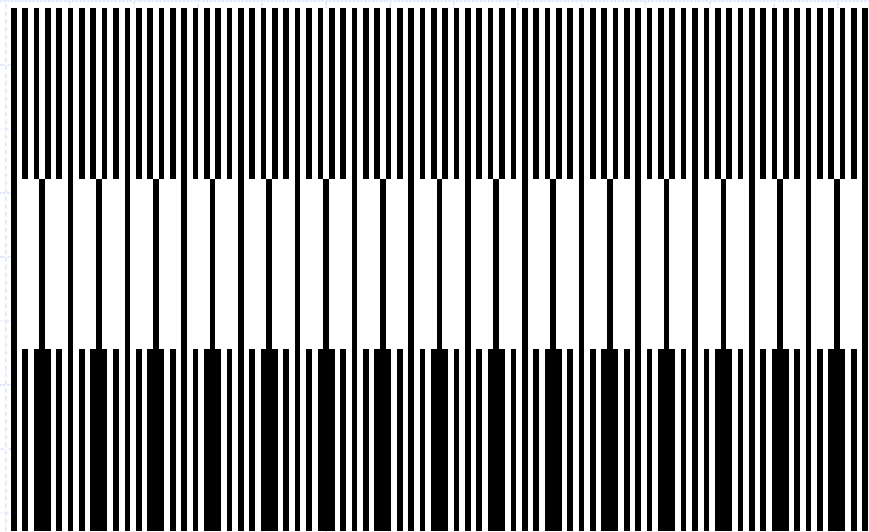
$$\tilde{A} \exp[i(kx - \omega t)] \text{ where } \tilde{A} \text{ is complex}$$

Again, the  $\sim$  is usually understood so the solution becomes

$$A \exp[i(kx - \omega t)]$$

# Waves

- How can we represent a (localized) particle by a (non-localized) wave?
  - Moire pattern



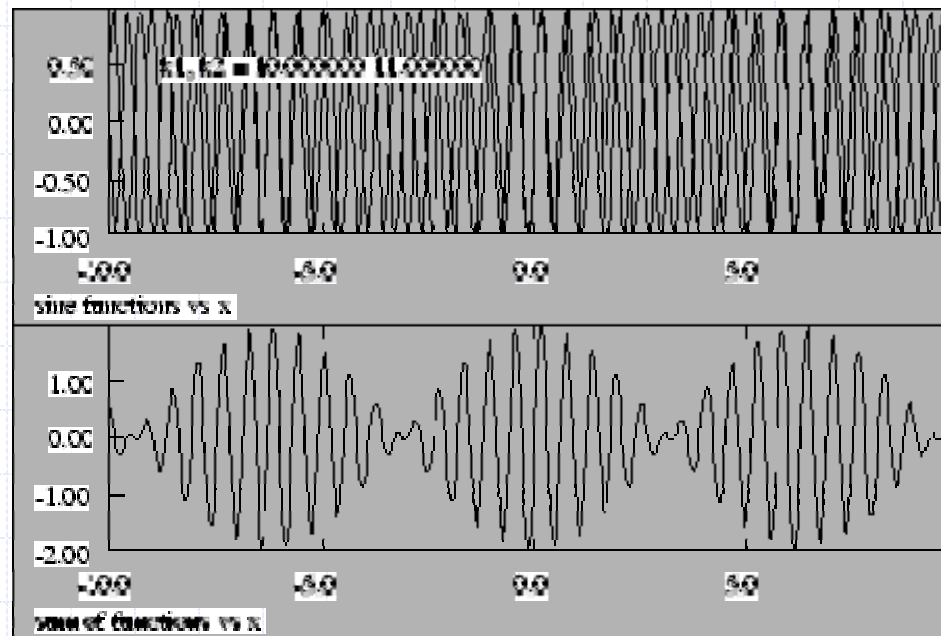
**Pattern 1**

**Pattern 2**

**1+2**

# Waves

- We can form a wave packet by summing waves of different wavelengths or frequencies



- Summing over an infinitely large number of waves between a range of  $k$  values will eliminate the auxiliary wave packets

# Waves

➤ Let's start by adding two waves together

$$\Psi(x, t) = \Psi_1(x, t) + \Psi_2(x, t)$$

$$\Psi(x, t) = A \cos(k_1 x - \omega_1 t) + A \cos(k_2 x - \omega_2 t)$$

$$\text{using } \cos A + \cos B = 2 \cos\left(\frac{A+B}{2}\right) \cos\left(\frac{A-B}{2}\right)$$

$$\Psi(x, t) = 2A \cos(k_{av} x - \omega_{av} t) \cos\left(\frac{\Delta k}{2} x - \frac{\Delta \omega}{2} t\right)$$

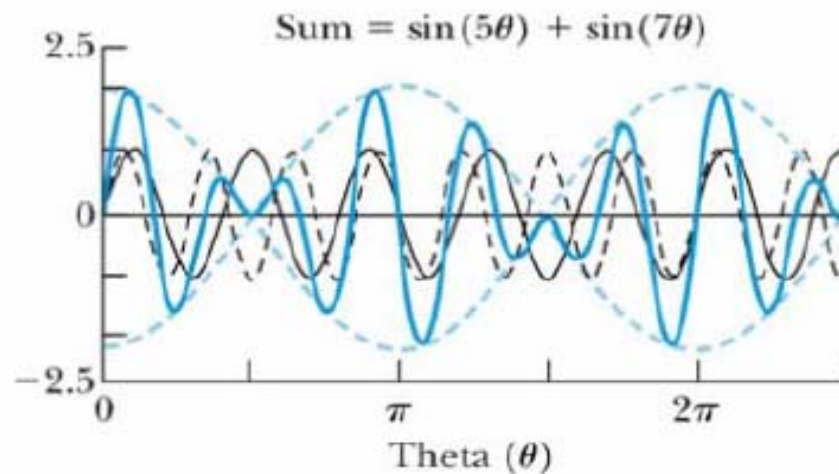
$$\text{where } k_{av} = \frac{k_1 + k_2}{2} \text{ and } \Delta k = k_1 - k_2$$

and similarly for  $\omega$

# Waves

- The first term is the wave and the second term is the envelope

$$\Psi(x, t) = 2A \cos(k_{av}x - \omega_{av}t) \cos\left(\frac{\Delta k}{2}x - \frac{\Delta\omega}{2}t\right)$$



# Waves

➤ The combined wave has phase velocity

$$v_{phase} = \frac{\omega_{av}}{k_{av}}$$

➤ The group (envelope) has group velocity

$$v_{group} = \frac{\Delta\omega}{\Delta k}$$

➤ More generally the group velocity is

$$v_{group} = \frac{d\omega(k)}{dk}$$

# Waves

- Recall our problem with the wave velocity being slower than the particle velocity
- Using the group velocity we find

$$v_{group} = \frac{d\omega}{dk} = \frac{d\hbar\omega}{d\hbar k} = \frac{dE}{dp}$$

$$E = \frac{p^2}{2m}$$

$$v_{group} = \frac{p}{m} = V$$

- Problem solved



# Waves

➤ The relation between  $\omega$  and  $k$  is called the dispersion relation

➤ Examples

■ Photon in vacuum

$$f = \frac{c}{\lambda}$$

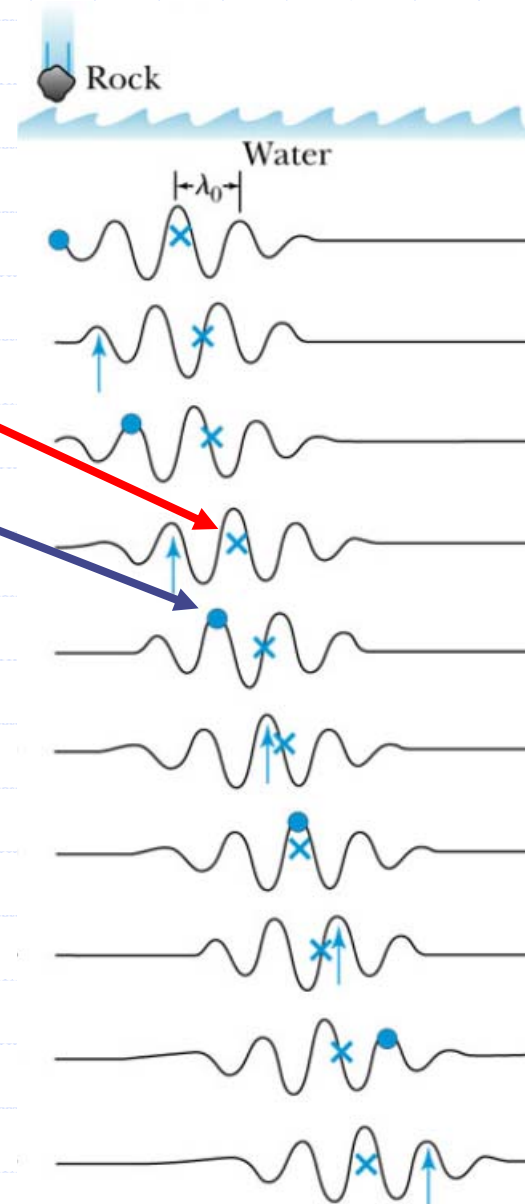
$$2\pi f = \frac{2\pi c}{\lambda}$$

$$\omega = kc$$

$$\frac{d\omega}{dk} = c$$

# Waves

➤ Phase and group velocity



# Waves

## ➤ Examples

- Photon in a medium

$$f = \frac{c}{n(\lambda)\lambda}$$

$$2\pi f = \frac{2\pi c}{n(\lambda)\lambda}$$

$$\omega = \frac{kc}{n(k)}$$

$$\frac{d\omega}{dk} = \frac{c}{n(k)} + kc \frac{d}{dk} \frac{1}{n(k)}$$

# Waves

## ➤ Examples

- de Broglie waves

$$E = hf = \hbar\omega$$

$$E = \frac{p^2}{2m} = \frac{(\hbar k)^2}{2m}$$

$$\omega = \frac{\hbar k^2}{2m}$$

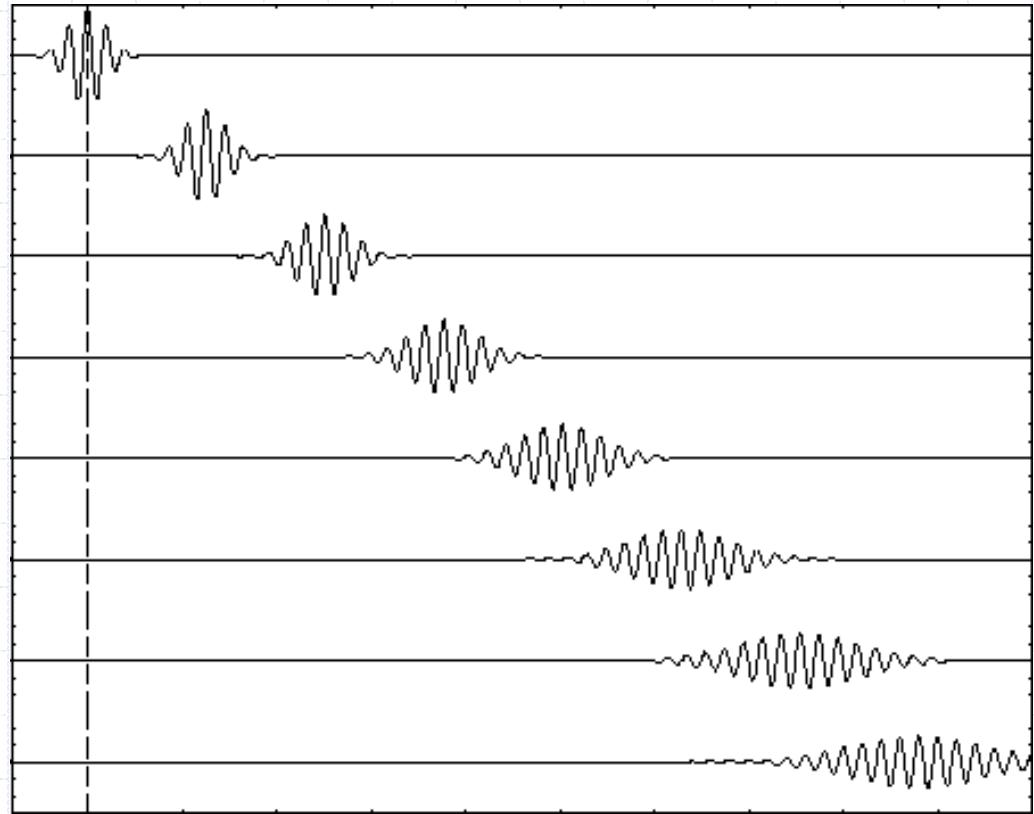
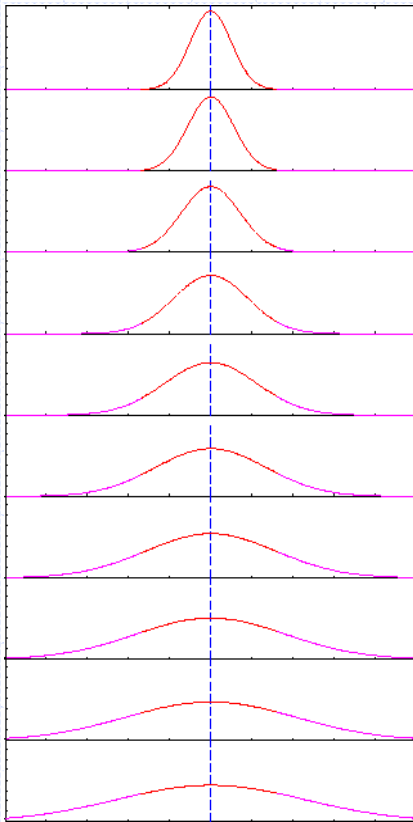
$$\frac{d\omega}{dk} = \frac{\hbar k}{m}$$

# Waves

- Because the group velocity =  $d\omega/dk$  depends on  $k$ , the wave packet will disperse with time
  - This is because each of the waves is moving at a slightly different velocity
  - Just as white light will be dispersed as it travels through a prism
- That's why  $\omega(k)$  is called a dispersion relation

# Waves

## ➤ Spreading of a wave packet with time



## ➤ Does this mean matter waves disperse with time?