

Cross Section

➤ Much of what we have learned about the internal structure of matter comes from scattering experiments (i.e. a microscope)

- Beam = electron, pion, proton, antiproton, alpha, photon, neutrino, strange particles, nuclei, ...
- Target = electron, proton, alpha, photon, nuclei, ...

➤ Fixed target experiments



➤ Colliding beam experiments



Cross Section

- Imagine shooting bullets at a circular target behind a curtain
- The bullets are spread uniformly

$$R = I\sigma$$

where

R is the hit frequency (bullets/sec)

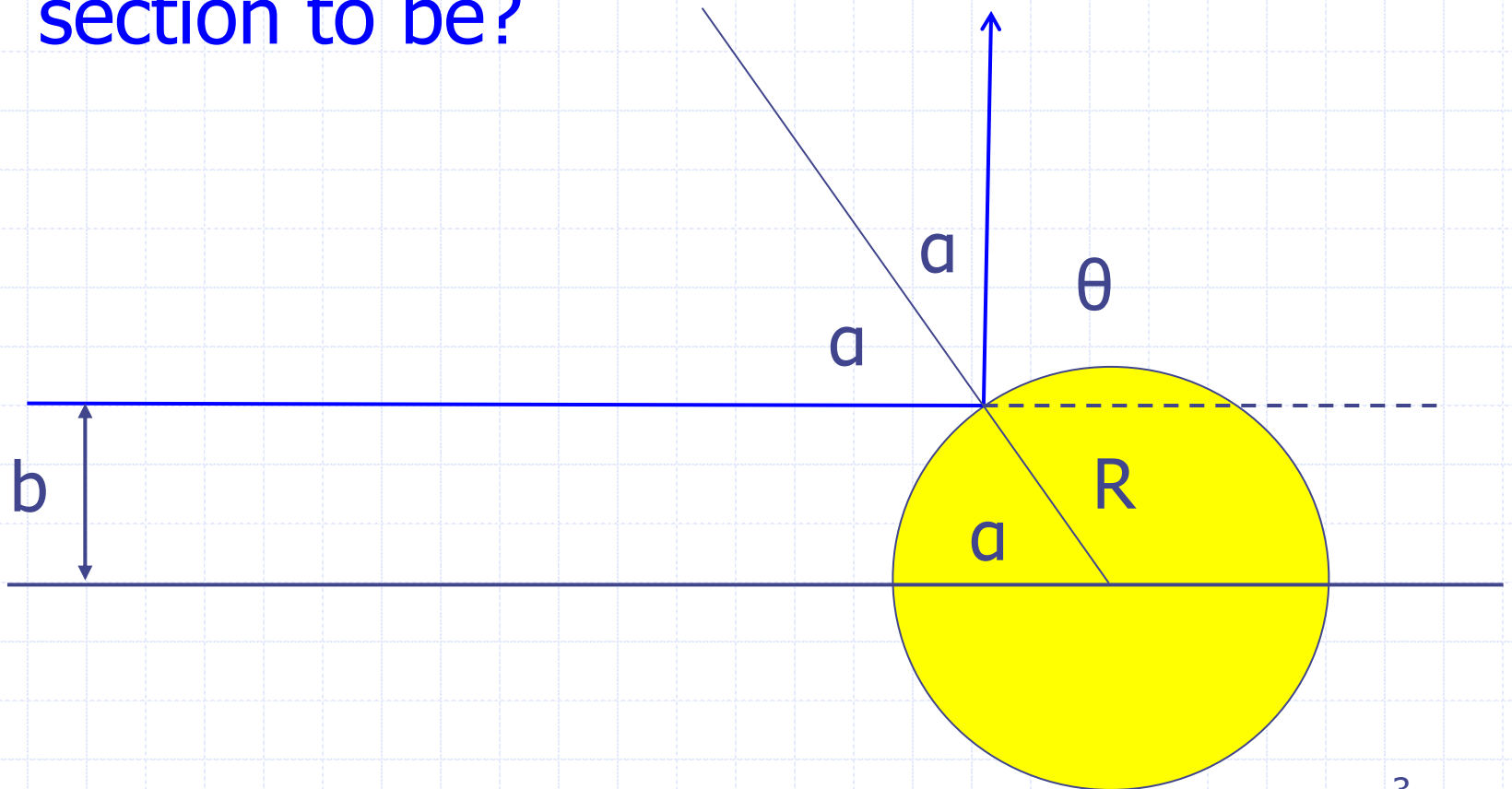
I is the incoming bullet intensity (bullets/m²/sec)

σ is the cross section (m²)

- The larger the target (σ) the more hits (scatters)

Cross Section

- Consider scattering from a hard sphere
- What would you expect the cross section to be?



Cross Section

➤ From the figure we see

$$b = R \sin \alpha \text{ and } 2\alpha + \theta = \pi$$

$$\text{then } \sin \alpha = \sin(\pi / 2 - \theta / 2) = \cos(\theta / 2)$$

$$\text{and } b = R \cos(\theta / 2)$$

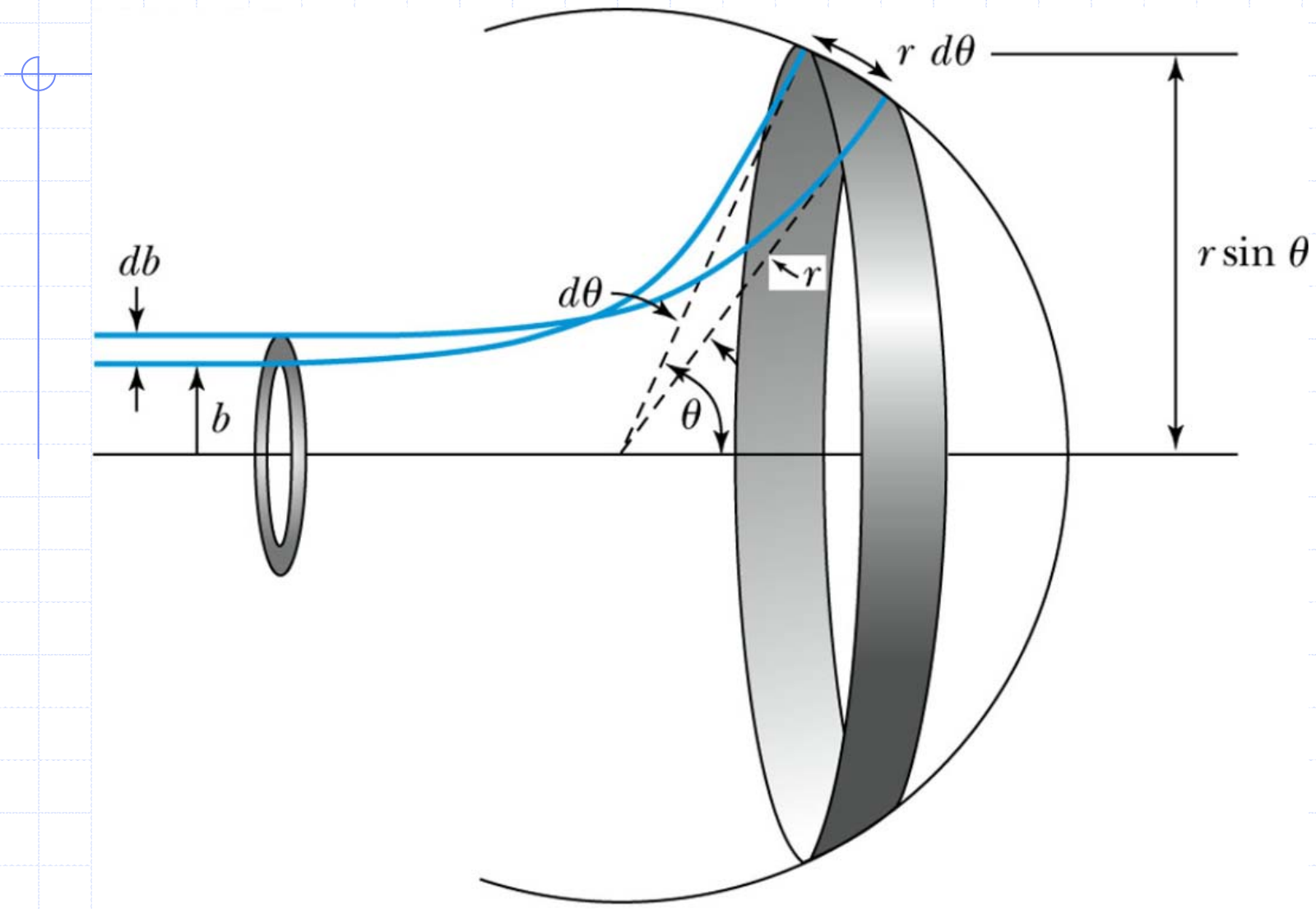
$$\text{or } \theta = 2 \cos^{-1}(b / R)$$

➤ This is the relation between b and θ for hard sphere scattering

Cross Section

- If a particle arrives with an impact parameter between b and $b+db$, it will emerge with a scattering angle between θ and $\theta+d\theta$
- If a particle arrives within an area of $d\sigma$, it will emerge into a solid angle $d\Omega$

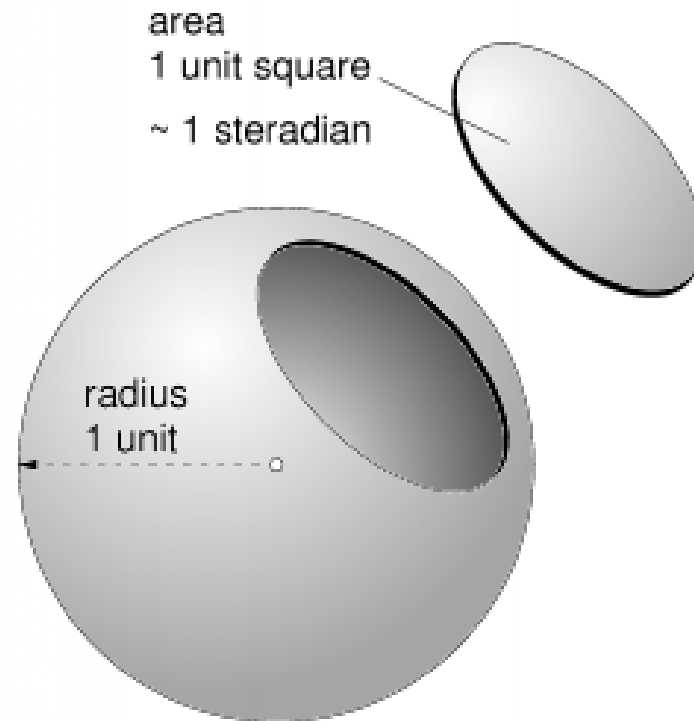
Cross Section



Cross Section

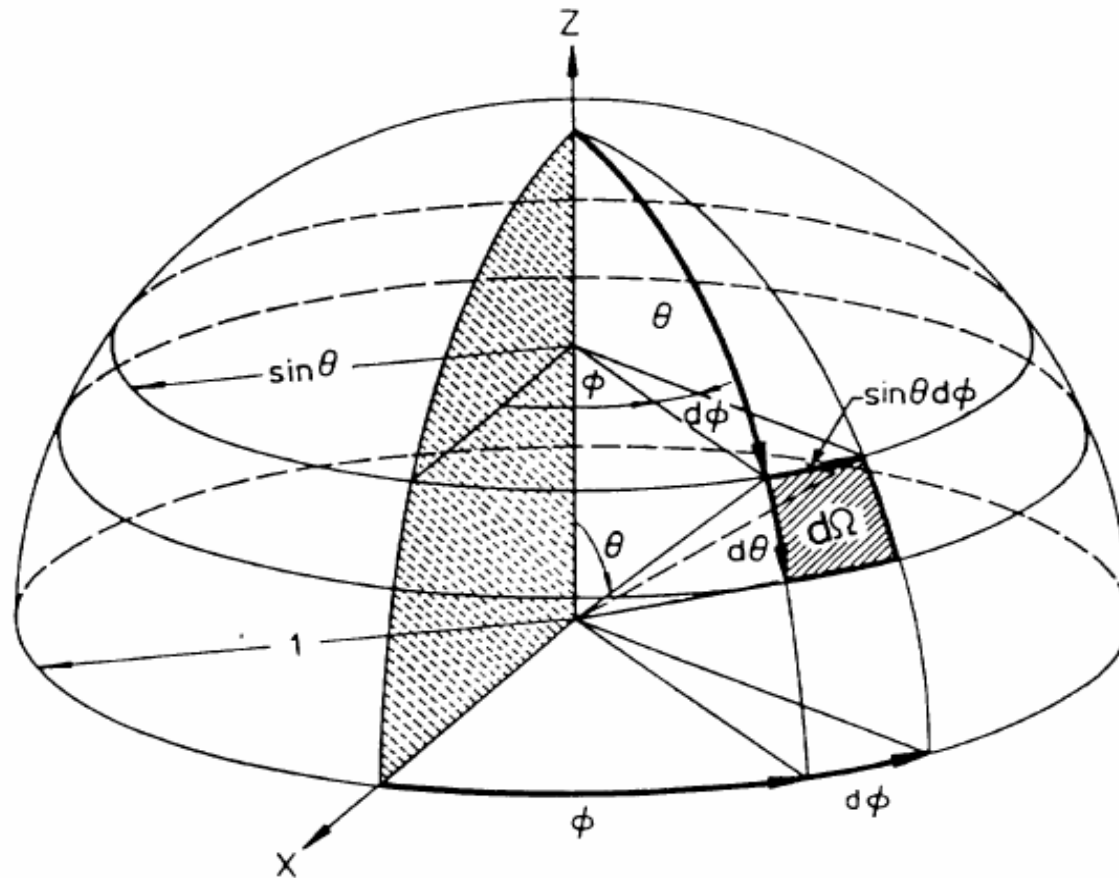
➤ $\Omega = A/r^2$

➤ So a sphere has 4π steradians (sr)



Cross Section

→ $d\Omega = dA/r^2 = \sin\theta d\theta d\phi$



Cross Section

➤ We have

$$d\sigma = \frac{d\sigma}{d\Omega} d\Omega$$

where $d\sigma = b db d\phi$

and $d\Omega = \sin \theta d\theta d\phi$

➤ And the proportionality constant $d\sigma/d\Omega$ is called the differential cross section

Cross Section

➤ Then we have

$$\frac{d\sigma}{d\Omega} = \left| \frac{b}{\sin \theta} \frac{db}{d\theta} \right|$$

➤ And for the hard sphere example

$$\frac{db}{d\theta} = -\frac{R}{2} \sin \frac{\theta}{2}$$

$$\frac{d\sigma}{d\Omega} = \frac{Rb \sin \frac{\theta}{2}}{2 \sin \theta} = \frac{R^2 \cos \frac{\theta}{2} \sin \frac{\theta}{2}}{2 \sin \theta} = \frac{R^2}{4}$$

Cross Section

➤ Finally

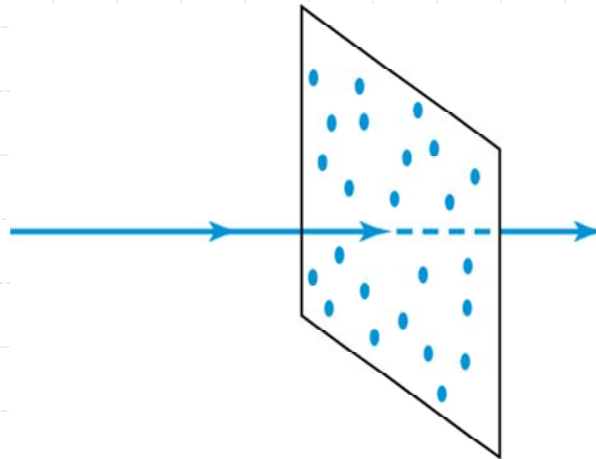
$$\sigma = \int d\sigma = \int \frac{d\sigma}{d\Omega} d\Omega = \int \frac{R^2}{4} d\Omega = \pi R^2$$

- This is just as we expect
- The cross section formalism developed here is the same for any type of scattering (Coulomb, nuclear, ...)

Cross Section

➤ The units of cross section are barns

- 1 barn (b) = 10^{-28}m^2
- The units are area. One can think of the cross section as the effective target area for collisions. We sometimes take $\sigma = \pi r^2$



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Cross Section

➤ One can find the scattering rate by

- $R = I_0 N_T \sigma$

- ◆ number/s = (number/s) (number nuclei/cm²) (cm²)

- $N_T = N_{Av} \rho l / At$

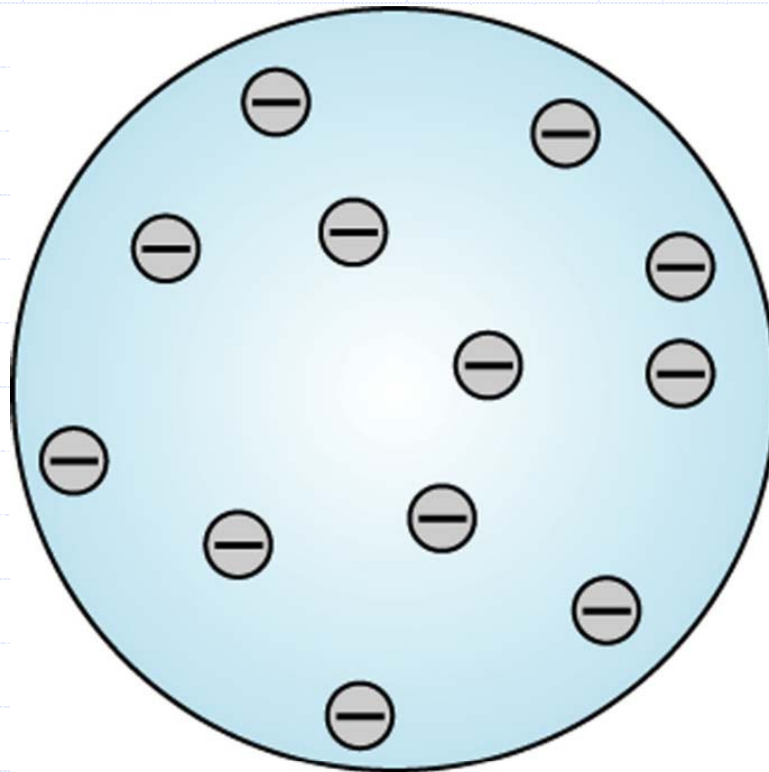
- ◆ nuclei/cm² = (Avagadro's number) (density) (length) / (atomic weight in g)

➤ And the scattering rate into dΩ is

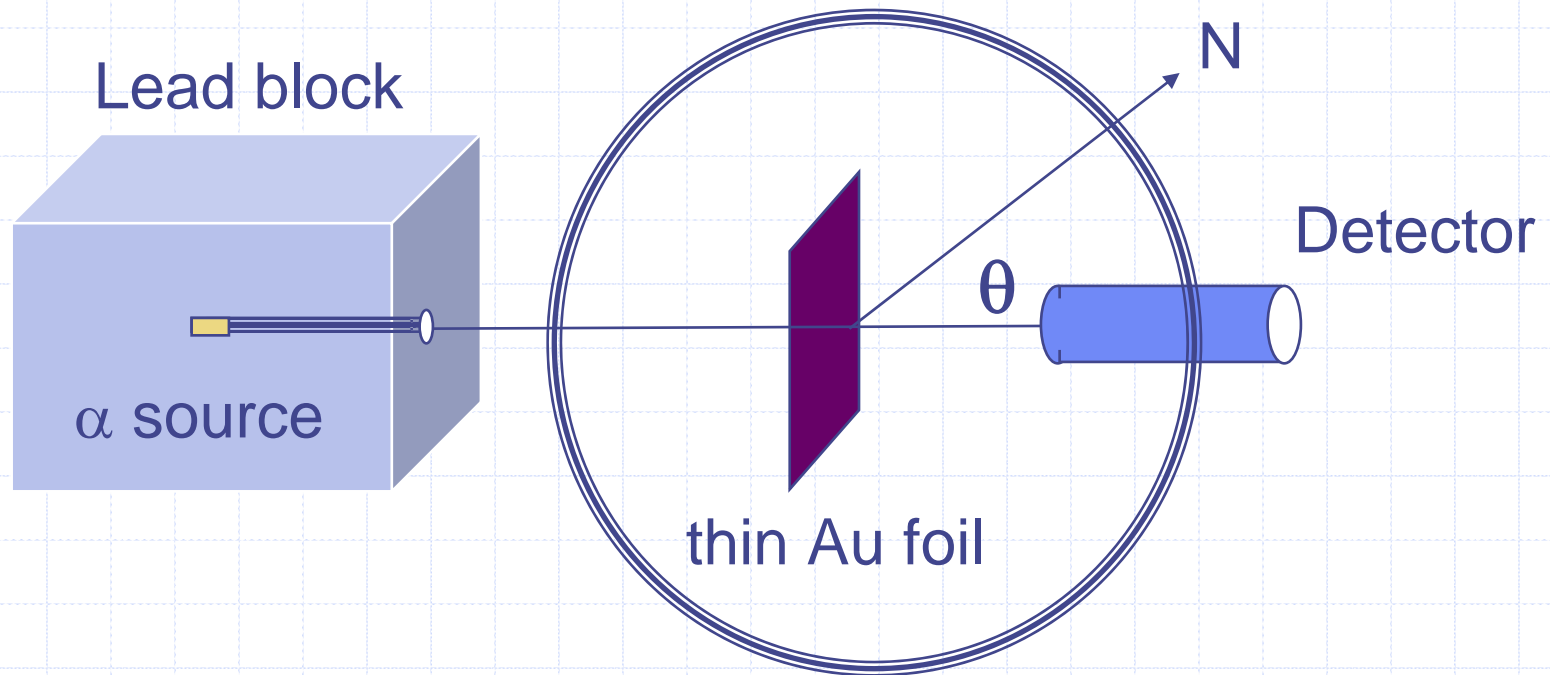
- $R(d\Omega) = I_0 N_T d\sigma/d\Omega$

Thomson's Plum Pudding Model

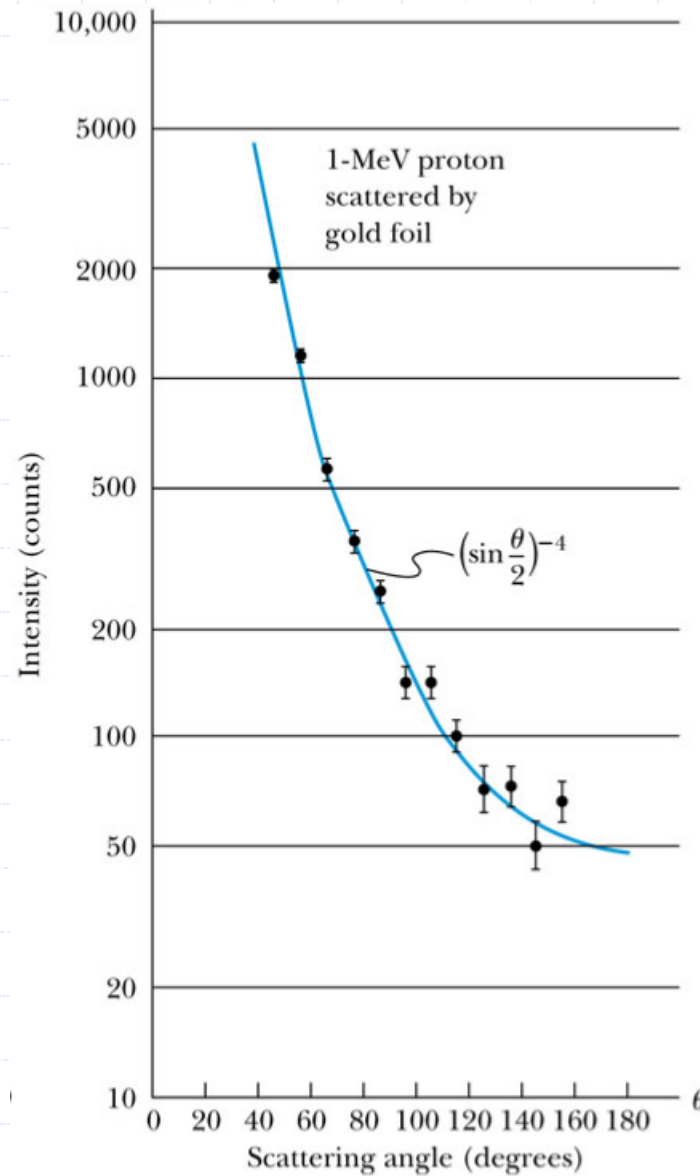
- Sphere diameter $\sim 10^{-10}\text{m}$
- Problems explaining line spectra and electrical stability



Rutherford Scattering



Rutherford Scattering



Rutherford Scattering

➤ Scattering from electrons?

- Recall from mechanics that a head-on collision from an alpha with momentum MV with an electron of mass m at rest gives a momentum change of $2mV$ for the electron
- The momentum change of the alpha is approximately $2mV=2MV/8000=MV/4000$
- An upper limit for the deflection can be found by assuming the momentum change is perpendicular to the momentum

$$\frac{\Delta p}{p} \approx \tan \theta \approx \theta \approx \frac{1}{4000} \approx 0.01^\circ$$

- Even multiple scattering from many (thousands) of electrons does not lead to scattering angles of much more than $\sim 1^\circ$

Rutherford Scattering

➤ Scattering from the plum pudding?

- The maximum Coulomb force on an alpha of charge q a distance R from a charge Q is
- The momentum change is
- Again taking the change in momentum to be perpendicular to the momentum

$$F = \frac{kqQ}{R^2}$$

$$\Delta p \approx F\Delta t = \frac{kqQ}{R^2} \frac{2R}{V}$$

$$\tan \theta \approx \theta \approx \frac{\Delta p}{p} = \frac{kqQ}{R \frac{1}{2} MV^2}$$

$$\theta \approx \frac{(2)(79)(14.4eV - A)}{(1A)(5 \times 10^6 eV)} = 0.026^\circ$$

Rutherford Scattering

➤ A convenient form for ke^2

$$\begin{aligned} ke^2 &= \left(9 \times 10^9 N \frac{m^2}{C^2} \right) (1.6 \times 10^{-19} C)^2 \\ &= (9 \times 10^9) (1.6 \times 10^{-19})^2 J - m \times \frac{1eV}{1.6 \times 10^{-19} J} \\ &= 14.4 \times 10^{-10} eV - m \\ &= 14.4 eV - Angstroms(A) \end{aligned}$$

Rutherford Scattering

- The complete calculation of the Rutherford scattering cross section can be found in Thornton and Rex (section 4.2)
- Instead I will do a poor man's calculation to show the idea
 - Assume a charged particle Z_1e Coulomb scatters from a nucleus of charge Z_2e
 - Assume the nucleus is infinitely massive

Rutherford Scattering

Electric force $F = k \frac{Z_1 Z_2 e^2}{b^2}$; $k = \frac{1}{4\pi\epsilon_0}$

Interaction time $\Delta t = \frac{2b}{v}$

Impulse $I = \Delta p_T \approx F\Delta t = k \frac{2Z_1 Z_2 e^2}{bv}$

Scattering angle $\theta \approx \frac{\Delta p_T}{p} = k \frac{2Z_1 Z_2 e^2}{pbv}$

So $b = k \frac{2Z_1 Z_2 e^2}{pv\theta}$

Aside, the exact relation is $b = k \frac{Z_1 Z_2 e^2}{pv} \cot \frac{\theta}{2}$

And note $\cot \frac{\theta}{2} \approx \left(\frac{\theta}{2}\right)^{-1} - \frac{1}{3} \frac{\theta}{2} - \dots$

Rutherford Scattering

Then $\frac{db}{d\theta} = -k \frac{2Z_1Z_2e^2}{pv\theta^2}$

Recalling our earlier result $\frac{d\sigma}{d\Omega} = \left| \frac{b}{\sin\theta} \frac{db}{d\theta} \right|$

And using $\sin\theta \approx \theta$ and $b = k \frac{2Z_1Z_2e^2}{pv\theta}$

We arrive at

$$\frac{d\sigma}{d\Omega} = k^2 \frac{4Z_1^2Z_2^2e^4}{p^2v^2} \frac{1}{\theta^4}$$

That you can compare to the exact Rutherford scattering cross section

$$\frac{d\sigma}{d\Omega} = k^2 \frac{Z_1^2Z_2^2e^4}{p^2v^2} \frac{1}{4\sin^4\left(\frac{\theta}{2}\right)}$$

Rutherford Scattering

➤ Rutherford scattering differential cross section

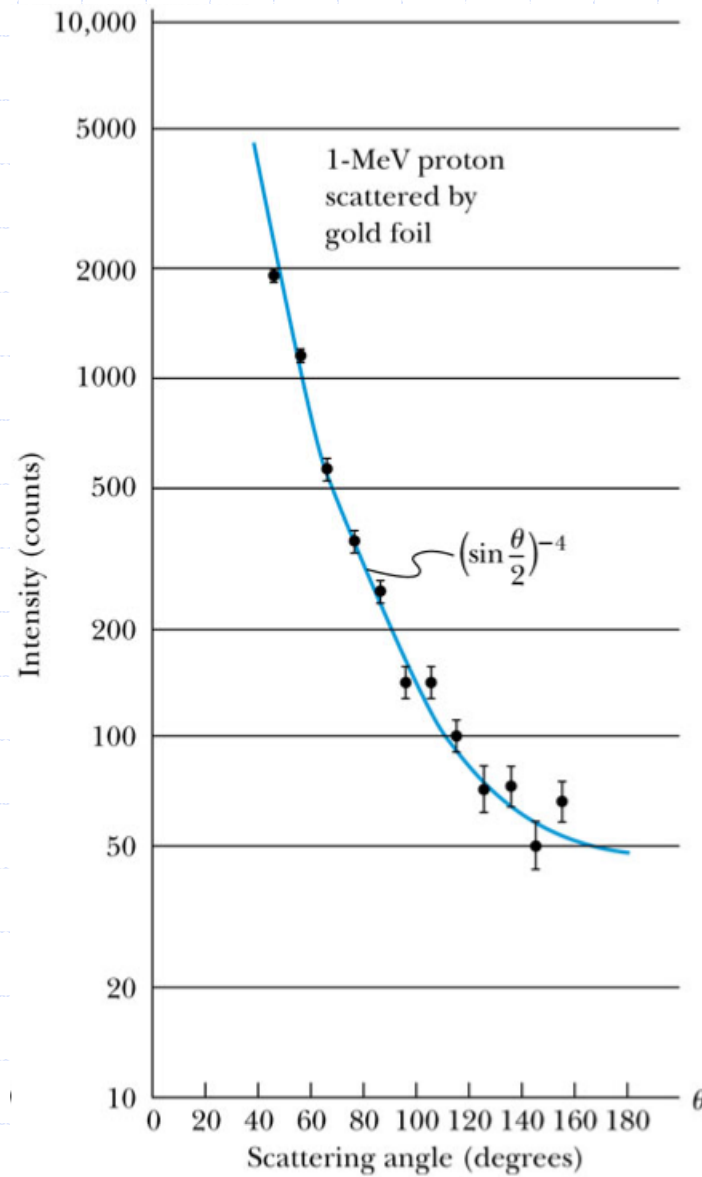
$$\frac{d\sigma}{d\Omega} = k^2 \frac{Z_1^2 Z_2^2 e^4}{p^2 v^2} \frac{1}{4 \sin^4\left(\frac{\theta}{2}\right)}$$

➤ The important features are

- $1/\sin^4(\theta/2)$ dependence
- Z_1^2, Z_2^2 dependence
- $1/T^2$ dependence

➤ All observed experimentally

Rutherford Scattering



Rutherford Scattering

➤ Consider scattering from the plum pudding

- We estimated the scattering angle to be $\theta \sim 0.026^\circ$
- Invoking the central limit theorem-the sum of a large number of independent variables approaches a Gaussian distribution

$$N = N_0 e^{-\left(\frac{\theta}{\theta_m}\right)^2}$$

➤ The expected number scattered through 90° or more is

$$N_{90} = N_0 e^{-\left(\frac{90}{1}\right)^2} = N_0 e^{-8100} \approx N_0 10^{-3500} \approx 0$$

Rutherford Scattering

➤ The relation between impact and scattering angle is

$$b = \frac{Z_1 Z_2 e^2}{8\pi\epsilon_0 T} \cot \frac{\theta}{2}$$

➤ The scattering is deterministic

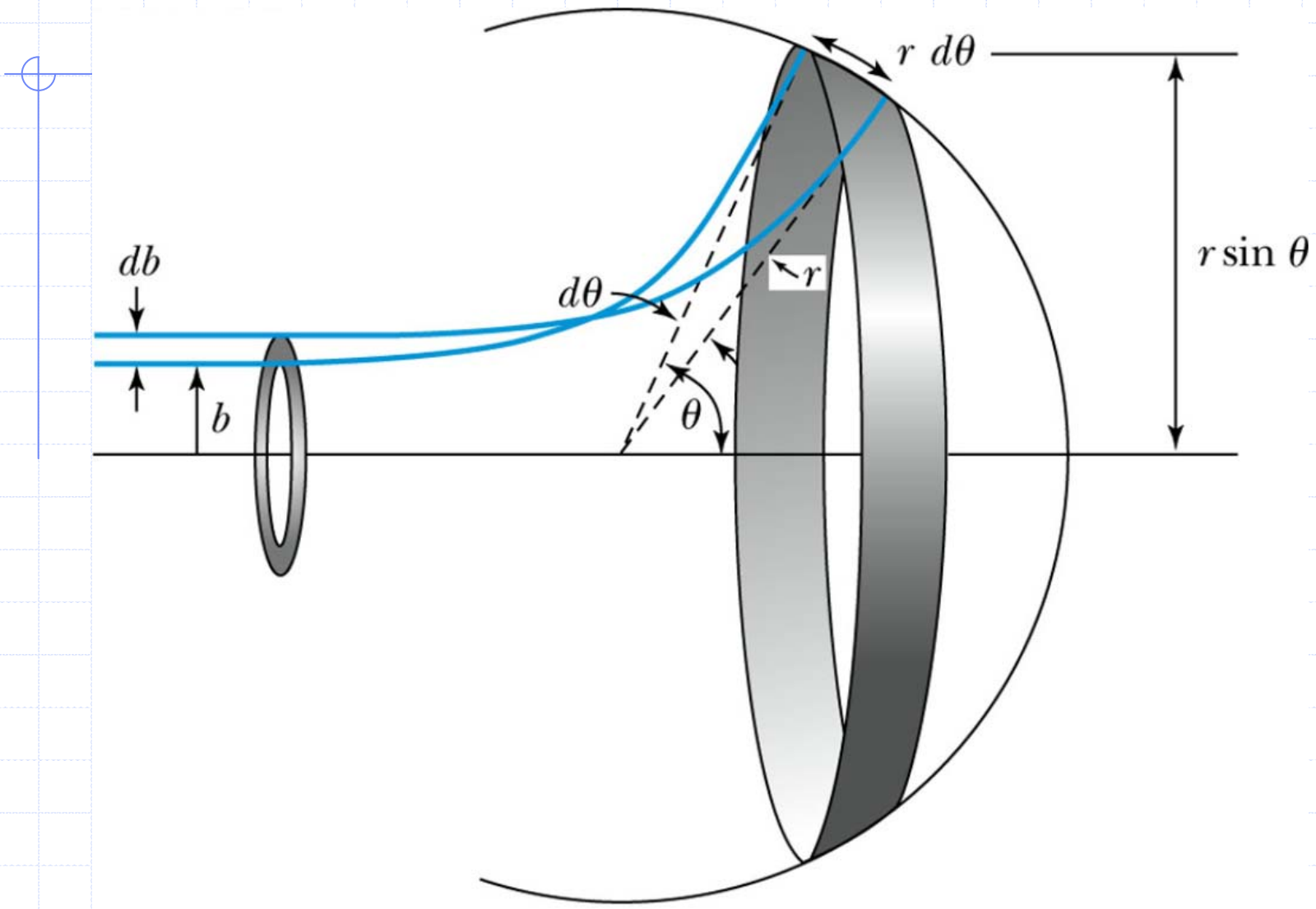
- This means all alpha particles with impact parameter $< b$ will scatter into angles $> \theta$
- So the number of alpha particles scattered into angles $> \theta$ is

$$\pi b^2 \times N_T$$

where N_T is the number of target nuclei/area

$$N_T = N_{Av} \rho t / A$$

Cross Section



Rutherford Scattering

➤ What fraction of particles is scattered through angles $> 90^\circ$

- 7.7 MeV alpha incident on an Au target of thickness $t=10^{-6}\text{m}$

$$N_T = \frac{(6.02 \times 10^{23}) \left(19.3 \frac{\text{g}}{\text{cm}^3} \right) (10^{-6} \text{m})}{197 \frac{\text{g}}{\text{mol}}} = 5.9 \times 10^{22} \frac{\text{nuclei}}{\text{m}^2}$$

$$f = \pi b^2 \times N_T$$

$$f = \pi (5.9 \times 10^{22}) \left(\frac{(79)(2)(14.4 \text{eVA})}{(2)(7.7 \text{MeV})} \right)^2 \cot^2 45^\circ$$

$$f = 4 \times 10^{-5}$$

Rutherford Scattering

- How large is the nucleus?
- Consider a head-on collision (180° scattering) of a 5 MeV alpha with a gold nucleus
- The distance of closest approach is found from

$$\frac{kqQ}{d} = \frac{1}{2}mV^2$$

$$d = \frac{kqQ}{\frac{1}{2}mV^2}$$

$$d = \frac{(2)(79)14.4eV - A}{5MeV} = 4.5 \times 10^{-14} m = 45F$$

Rutherford Scattering

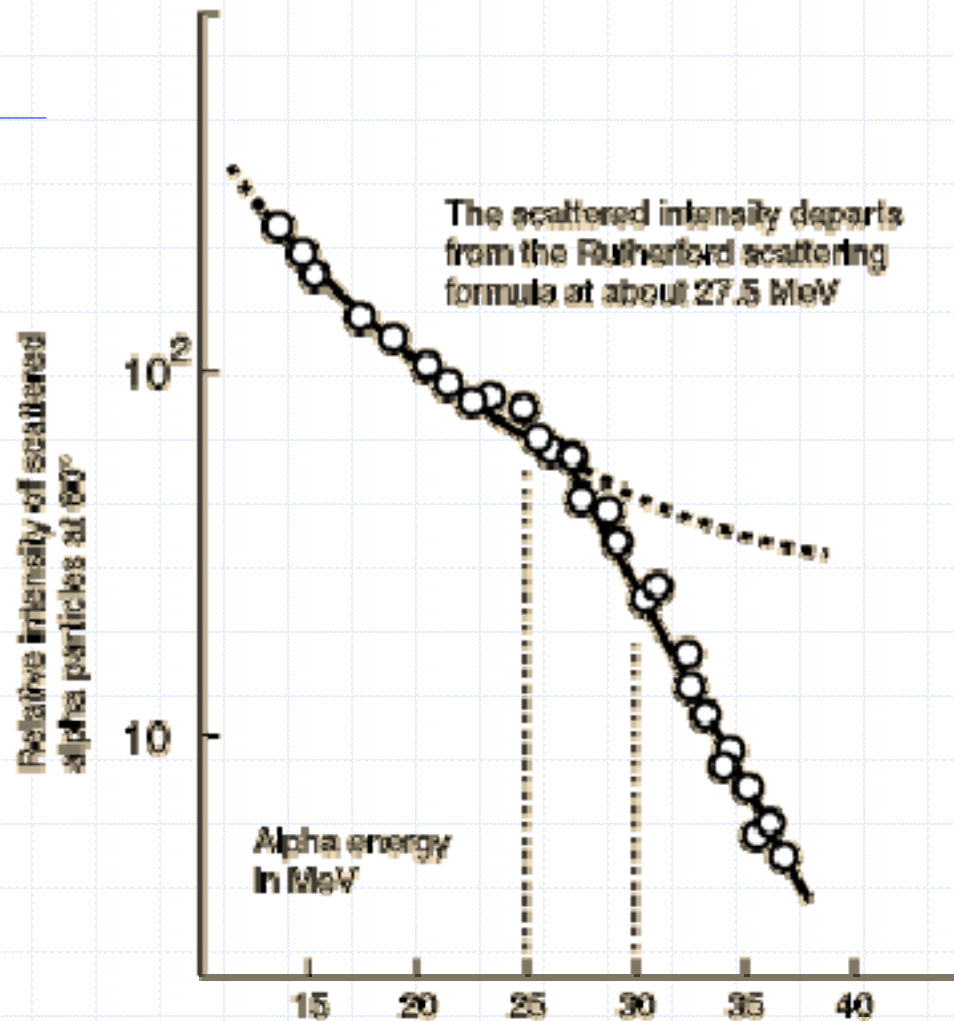
➤ The total cross section for Rutherford scattering is infinite!

$$\sigma = \int \frac{d\sigma}{d\Omega} d\Omega$$

$$\sigma = k^2 \frac{Z_1^2 Z_2^2 e^4}{4p^2 v^2} \int_0^\pi \frac{1}{\sin^4\left(\frac{\theta}{2}\right)} \sin\theta d\theta d\varphi$$

$$\sigma = \infty$$

Rutherford Scattering



➤ Welcome to nuclear physics